

# Orbital-specific mapping of the ligand exchange dynamics of Fe(CO)<sub>5</sub> in solution

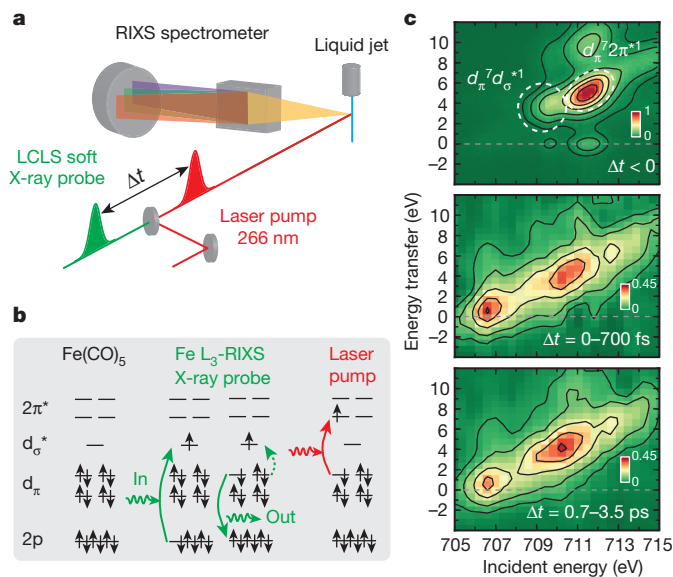
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Transition-metal complexes have long attracted interest for fundamental chemical reactivity studies and possible use in solar energy conversion<sup>1,2</sup>. Electronic excitation, ligand loss from the metal centre, or a combination of both, creates changes in charge and spin density at the metal site<sup>3–11</sup> that need to be controlled to optimize complexes for photocatalytic hydrogen production<sup>8</sup> and selective carbon–hydrogen bond activation<sup>9–11</sup>. An understanding at the molecular level of how transition-metal complexes catalyse reactions, and in particular of the role of the short-lived and reactive intermediate states involved, will be critical for such optimization. However, suitable methods for detailed characterization of electronic excited states have been lacking. Here we show, with the use of X-ray laser-based femtosecond-resolution spectroscopy and advanced quantum chemical theory to probe the reaction dynamics of the benchmark transition-metal complex Fe(CO)<sub>5</sub> in solution, that the photo-induced removal of CO generates the 16-electron Fe(CO)<sub>4</sub> species, a homogeneous catalyst<sup>12,13</sup> with an electron deficiency at the Fe centre<sup>14,15</sup>, in a hitherto unreported excited singlet state that either converts to the triplet ground state or combines with a CO or solvent molecule to regenerate a penta-coordinated Fe species on a sub-picosecond timescale. This finding, which resolves the debate about the relative importance of different spin channels in the photochemistry of Fe(CO)<sub>5</sub> (refs 4, 16–20), was made possible by the ability of femtosecond X-ray spectroscopy to probe frontier-orbital interactions with atom specificity. We expect the method to be broadly applicable in the chemical sciences, and to complement approaches that probe structural dynamics in ultrafast processes.

In our experimental setup (Fig. 1a), the valence electronic structure of Fe(CO)<sub>5</sub> is probed with femtosecond-resolution resonant inelastic X-ray scattering (RIXS) at the Fe L<sub>3</sub>-edge (Fe L<sub>3</sub>-RIXS, illustrated in Fig. 1b). The frontier orbitals of ironpentacarbonyl, Fe(CO)<sub>5</sub>, and its photofragments are the Fe-centred  $d_{\pi}$  and  $d_{\sigma}^*$  orbitals. With an incident photon energy of 710 eV to select the lowest-energy X-ray resonance corresponding to  $2p \rightarrow \text{LUMO}(d_{\sigma}^*)$  (where LUMO is the lowest unoccupied molecular orbital) excitations and scattering inelastically to the valence-excited ligand-field states with  $d_{\pi}^7 d_{\sigma}^{*1}$  configuration, we effectively probe  $d_{\pi} \rightarrow d_{\sigma}^*$  transitions (note that the single-electron orbital-based assignments can be applied at the level that the system is studied here; see Supplementary Information). The energies of these transitions equal the measured energy transfers (that is, the difference between incident and scattered photon energies indicated by ‘in’ and ‘out’ in Fig. 1b), and directly reflect the changes in chemical bonding and ligand coordination. The intensities of the transitions in Fe(CO)<sub>5</sub> are marked

in Fig. 1c (top). The main intensity maximum involves  $2p \rightarrow 2\pi^*$  excitations at 711.5 eV with excitation to the ligand-centred  $2\pi^*$  orbitals and inelastic scattering to  $d_{\pi}^7 d_{\sigma}^{*1}$  charge-transfer states (Fig. 1c) and is not further analysed.

The unsaturated carbonyl Fe(CO)<sub>4</sub> was generated in ethanol (EtOH) solution by the photodissociation of Fe(CO)<sub>5</sub> with optical (266 nm)



**Figure 1 | Scheme and results of the experiment.** **a**, Scheme with optical-laser pump and soft X-ray probe after the pump–probe time delay  $\Delta t$ . The intensity of RIXS is measured at the Fe L<sub>3</sub>-absorption edge with a dispersive grating spectrometer. **b**, Electron configuration of ground-state Fe(CO)<sub>5</sub> with single-electron transitions of X-ray probe and laser-pump processes (orbital assignments according to Fe  $2p$  and  $3d$  or ligand  $2\pi$  character and according to symmetry along the Fe–CO bonds; the asterisk marks antibonding orbitals). RIXS at the Fe L<sub>3</sub>-absorption edge with  $2p \rightarrow d_{\sigma}^*$  excitation involves scattering to final  $d_{\pi}^7 d_{\sigma}^{*1}$  ligand-field excited states. Optical  $d_{\pi} \rightarrow 2\pi^*$  excitation triggers dissociation. **c**, Measured Fe L<sub>3</sub>-RIXS intensities (encoded in colour) versus energy transfer and incident photon energy. Top: ground-state Fe(CO)<sub>5</sub> (negative delays, probe before pump). Middle and bottom: difference intensities for delay intervals of 0–700 fs and 0.7–3.5 ps, respectively, isolating transients by subtracting scaled intensities of unpumped Fe(CO)<sub>5</sub> from the measured intensities (scaling factor 0.9). For details of the experiment and a deduction of the scaling factor see Supplementary Information.

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femtosecond laser pulses in less than 100 fs. Our experiment consisted of recording Fe  $L_3$ -RIXS intensities versus energy transfer while scanning incident photon energy and pump-probe time delay with a time resolution of 300 fs. The observed bimodal spectral distribution shows different intensities for different delays (Fig. 1c, middle and bottom), reflecting changes in  $2p \rightarrow$ LUMO resonance energies within the range 706.5–710 eV and changes in  $d_\pi \rightarrow d_\sigma^*$  transition energies within the

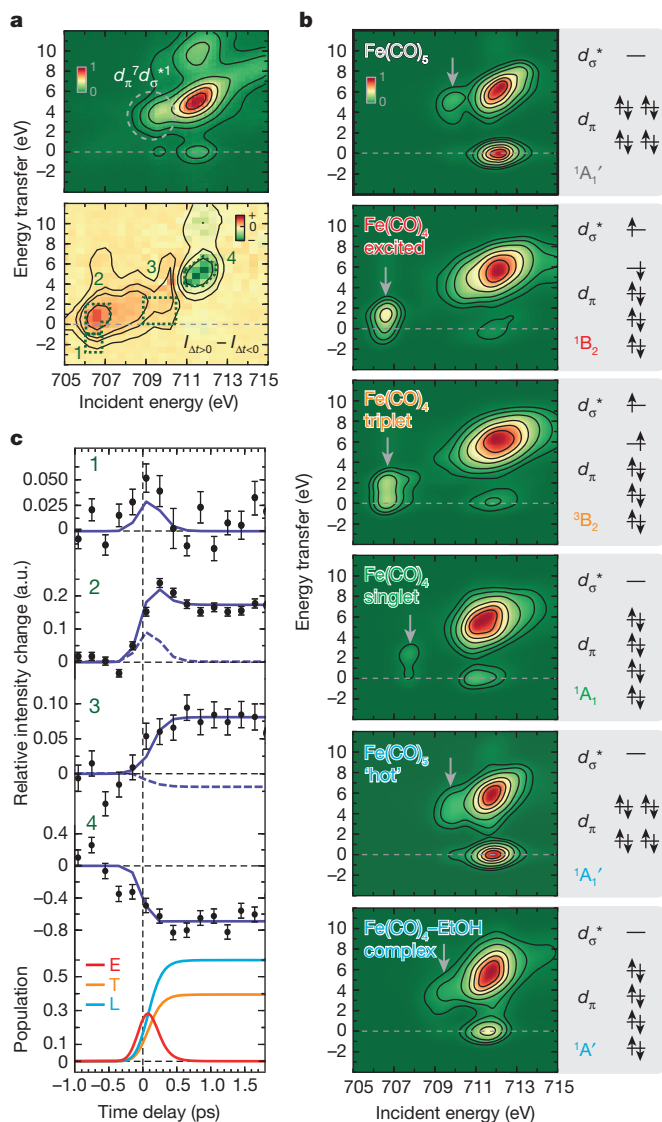
range –1 to 6 eV. These changes quantify the changes in the frontier-orbital interactions caused by changes in the ligand coordination when going from  $\text{Fe}(\text{CO})_5$  to  $\text{Fe}(\text{CO})_4$  and during the subsequent excited-state dynamics. Ligand dissociation is expected to create a ‘localized hole on the metal’<sup>15</sup> with a concomitant decrease in the  $d_\pi$ – $d_\sigma^*$  splitting (see the molecular-orbital diagram in Supplementary Information). This manifests itself in the Fe  $L_3$ -RIXS spectra at time delays of 0–700 fs (Fig. 1c, middle) as a new  $2p \rightarrow$ LUMO resonance at 706.5 eV and as the maximum of the  $d_\pi \rightarrow d_\sigma^*$  transitions shifted to lower energies by –4 eV relative to  $\text{Fe}(\text{CO})_5$ . Coordinative saturation through ligation with CO or EtOH restores the  $d_\pi$ – $d_\sigma^*$  splitting, mostly as a result of  $\sigma$ -bonding between  $\text{Fe}(\text{CO})_4$  and CO or EtOH. This could explain the occurrence of  $2p \rightarrow$ LUMO and  $d_\pi \rightarrow d_\sigma^*$  transition energies comparable to  $\text{Fe}(\text{CO})_5$  at late delays of 0.7–3.5 ps (at 709.5 and 3 eV; Fig. 1c, bottom).

To substantiate this and to quantitatively analyse the time-resolved data in Fig. 2a, we performed *ab initio* Fe  $L_3$ -RIXS calculations for selected structures. The calculated spectra of the three lowest electronic states of  $\text{Fe}(\text{CO})_4$ , of the lowest states of  $\text{Fe}(\text{CO})_4$ –EtOH complexes and of  $\text{Fe}(\text{CO})_5$  in optimized and distorted geometries account for all experimental features. Figure 2b shows the spectra and electronic configurations corresponding to the excited singlet-state  $\text{Fe}(\text{CO})_4$  ( $d_\pi^7 d_\sigma^{*1}$ ,  $^1B_2$ ), triplet-state  $\text{Fe}(\text{CO})_4$  ( $d_\pi^7 d_\sigma^{*1}$ ,  $^3B_2$ ), singlet-state  $\text{Fe}(\text{CO})_4$  ( $d_\pi^8 d_\sigma^{*0}$ ,  $^1A_1$ ), ‘hot’ singlet  $\text{Fe}(\text{CO})_5$  ( $d_\pi^8 d_\sigma^{*0}$ ,  $^1A_1'$ ), as represented by structures with distorted geometries compared to the optimized one) and singlet complexes with the solvent  $\text{Fe}(\text{CO})_4$ –EtOH ( $d_\pi^8 d_\sigma^{*0}$ ,  $^1A'$ ).

The most informative spectral regions in our data, labelled 1–4 in Fig. 2a, overlap maximally with the calculated spectral features best able to identify and distinguish the respective intermediate species. The intensities at negative transfers in region 1 result from outgoing X-rays with higher energy than the incoming X-rays and can only result when the X-rays scatter inelastically off  $\text{Fe}(\text{CO})_4$  fragments in  $d_\pi$  and  $d_\sigma^*$  electronic excited states. Region 2 is dominated by contributions of excited and triplet  $\text{Fe}(\text{CO})_4$ . We emphasize that RIXS gives unique chemical resolution, because integrating over the energy transfer and measuring only time-dependent changes in X-ray absorption would prevent us from distinguishing the dynamics of the species assigned to regions 1 and 2. Region 3 identifies the dynamics of ligated  $\text{Fe}(\text{CO})_4$  species; that is,  $\text{Fe}(\text{CO})_5$  and  $\text{Fe}(\text{CO})_4$ –EtOH. Region 4 corresponds to the  $2p \rightarrow 2\pi^*$  X-ray resonances with  $d_\pi \rightarrow 2\pi^*$  RIXS transitions in most of the calculated species, and seems dominated by the depletion of  $\text{Fe}(\text{CO})_5$ .

The temporal evolution of the Fe  $L_3$ -RIXS intensities measured in regions 1–4 is plotted in Fig. 2c, together with the result of a kinetic model that simultaneously fits the sum of the calculated excited-state singlet, triplet and ligated  $\text{Fe}(\text{CO})_4$  spectra in each region to the measured data (see Supplementary Information for details of the kinetic model and the contribution of singlet  $\text{Fe}(\text{CO})_4$ ). This fitting procedure indicates the appearance of the excited singlet state of  $\text{Fe}(\text{CO})_4$  ( $^1B_2$ ) within the time resolution of our experiment (which is insufficient to resolve the initial, ultraviolet-generated excited state of  $\text{Fe}(\text{CO})_5$ ). This allows us to assign unambiguously, within the single-electron orbital picture, the related RIXS intensities at negative energy transfers to  $2p \rightarrow d_\pi$  excitations in excited  $\text{Fe}(\text{CO})_4$  ( $d_\pi^7 d_\sigma^{*1}$ ,  $^1B_2$ ) with predominant inelastic scattering to states with  $d_\pi^8 d_\sigma^{*0}$  configuration. As is apparent from the molecular orbital diagram of excited  $\text{Fe}(\text{CO})_4$  in Fig. 2b, these transitions entail negative energy transfer because the incident photon energy is smaller than the scattered photon energy. The detection and characterization of electronic excited states free from background by non-excited states, enabled by RIXS at negative energy transfers, provides a powerful approach to studying the electronic excited states of chemically active molecules.

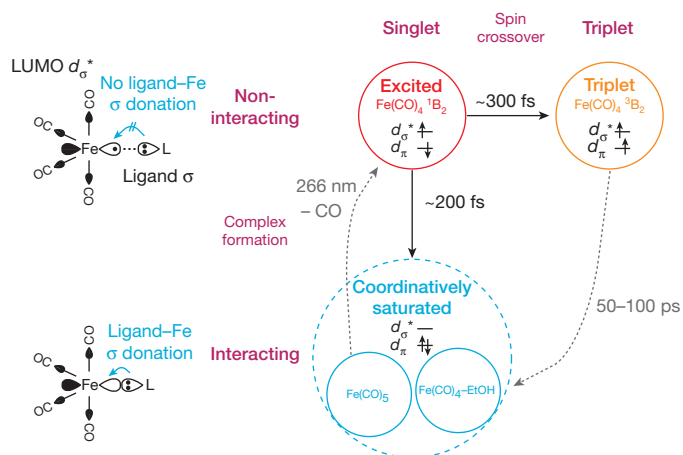
The decay of excited singlet-state  $\text{Fe}(\text{CO})_4$  ( $^1B_2$ ) coincides with the rise of the triplet  $\text{Fe}(\text{CO})_4$  ground state in solution ( $^3B_2$ ), for which our model gives a time constant of  $300 \pm 100$  fs. Within the experimental uncertainty, our data indicate the simultaneous rise of coordinatively saturated ‘hot’  $\text{Fe}(\text{CO})_5$  arising from geminate recombination with CO and of  $\text{Fe}(\text{CO})_4$ –EtOH arising from complexation with solvent molecules



**Figure 2 | Fe-specific changes in the electronic structure of  $\text{Fe}(\text{CO})_4$  after femtosecond spin crossover and ligation.** a, Measured Fe  $L_3$ -RIXS of  $\text{Fe}(\text{CO})_5$  (top, as in Fig. 1c) and measured difference intensities (bottom, integrated intensities of all positive pump-probe delays minus integrated intensities of all negative delays). The numbers 1–4 mark energy-transfer/incident-photon energy regions for which the temporal evolutions of intensities are plotted in c. b, Calculated Fe  $L_3$ -RIXS intensities and electronic configurations of the given species ( $2p \rightarrow$ LUMO and  $d_\pi \rightarrow d_\sigma^*$  transitions marked by arrows; the LUMO can be  $d_\sigma^*$  or  $d_\pi$ , depending on the electron configuration). c, Plot of measured intensities in regions 1–4 against pump-probe delay (means  $\pm$  s.d.) with the best global fit of a kinetic model (solid lines) with extracted populations of excited (E), triplet (T) and ligated (L)  $\text{Fe}(\text{CO})_4$  (L is a sum of ‘hot’  $\text{Fe}(\text{CO})_5$  and  $\text{Fe}(\text{CO})_4$ –EtOH). The dashed lines in regions 2 and 3 represent alternative models without triplet and ligated  $\text{Fe}(\text{CO})_4$ , respectively. The measured signals stayed constant up to 3 ps. For details of the calculations, structures and energies of the species and how ligation in  $\text{Fe}(\text{CO})_4$ –EtOH can occur through the alkyl or hydroxyl group see Supplementary Information.

(fitted time constant  $200 \pm 100$  fs). The failure of kinetic models without triplet  $\text{Fe}(\text{CO})_4$  ( $^3\text{B}_2$ ) (dashed curve in Fig. 2c, region 2) or without ligated  $\text{Fe}(\text{CO})_4$  (dashed curve in Fig. 2c, region 3) justifies the use of three distinct photoproducts in the kinetic modelling and underlines the robustness of our detection of triplet  $\text{Fe}(\text{CO})_4$  ( $^3\text{B}_2$ ) in parallel with 'hot'  $\text{Fe}(\text{CO})_5$  and  $\text{Fe}(\text{CO})_4\text{-EtOH}$ . Because we cannot spectroscopically distinguish geminately recombined 'hot'  $\text{Fe}(\text{CO})_5$  from solvent-complexed  $\text{Fe}(\text{CO})_4\text{-EtOH}$ , their ratio in the kinetic model is fixed at 1:1, consistent with the measured quantum yield of 0.8 for solvent-separated  $\text{Fe}(\text{CO})_4$  and CO (ref. 21).

Figure 3 sketches the reaction pathways established in this study, with detection of the excited singlet-state  $\text{Fe}(\text{CO})_4$  ( $^1\text{B}_2$ ) confirming the suggestion<sup>16</sup> that the primary reaction steps in solution also involve the singlet pathway as seen in the gas phase<sup>17,18</sup>. The proposed relaxation of excited singlet  $\text{Fe}(\text{CO})_4$  ( $^1\text{B}_2$ ) to singlet  $\text{Fe}(\text{CO})_4$  ( $^1\text{A}_1$ ) through internal conversion<sup>17</sup> is consistent with our data (see Supplementary Information), but we also observe triplet  $\text{Fe}(\text{CO})_4$  ( $^3\text{B}_2$ ) that was previously seen in solution<sup>4,19</sup> and in rare-gas matrix<sup>16</sup> experiments. This triplet arises from a singlet state with a time constant of 300 fs, consolidating the notion<sup>6</sup> that sub-picosecond intersystem crossing seems to be common in the excited-state dynamics of transition-metal complexes<sup>7,22–24</sup>. The persistence of the triplet  $\text{Fe}(\text{CO})_4$  ( $^3\text{B}_2$ ) up to our maximum time delay of 3 ps is consistent with its undergoing a slow, spin-forbidden reaction with intersystem crossing to a solvent-complexed singlet state on a 50–100 ps timescale<sup>4,5,25</sup>. However, the observed branching on a sub-picosecond timescale into the competing and simultaneous reaction channels of spin crossover and ligation to form coordinatively saturated species introduces an efficient pathway circumventing this spin barrier. It also supports the idea that the high density of electronic excited states and the relatively large amount of excess energy available in the system determine the course of the excited-state dynamics, rather than spin selection rules alone<sup>5,6</sup>. Fast ligation could be facilitated along the singlet pathway, confirming the general notion that solvent-stabilized metal centres form fast<sup>3,4,11</sup>; this is also consistent with the observation of the unsaturated carbonyl  $\text{Cr}(\text{CO})_5$  forming a solvent complex in alcohol solution within 1.6 ps (ref. 26). An alternative proposal<sup>20</sup> for  $\text{Fe}(\text{CO})_5$  involves a concerted exchange of CO and EtOH on the timescale of ligand dissociation of 100–150 fs. This would also proceed along a singlet pathway and would be in agreement with our results,



**Figure 3 | Schematic reaction pathways of  $\text{Fe}(\text{CO})_4$  in EtOH.** Parallel evolution from excited singlet-state  $\text{Fe}(\text{CO})_4$  to triplet-state  $\text{Fe}(\text{CO})_4$  through spin crossover (rise of triplet with a time constant of  $300 \pm 100$  fs) and to coordinatively saturated 'hot' singlet  $\text{Fe}(\text{CO})_5$  through geminate recombination and  $\text{Fe}(\text{CO})_4\text{-EtOH}$  by means of solvent-complex formation (increase in ligated  $\text{Fe}(\text{CO})_4$  with time constant  $200 \pm 100$  fs). A triplet pathway to  $\text{Fe}(\text{CO})_4\text{-EtOH}$  complex formation within 50–100 ps is indicated in grey. Interaction of the  $d_{\sigma}^*$  LUMO of  $\text{Fe}(\text{CO})_4$  with a ligand  $\sigma$  orbital (different phases are shown in black and white) is shown in the sketches with ligand–Fe  $\sigma$  donation in the coordinatively saturated species.

because the temporal resolution of our measurements is not sufficient to distinguish between this concerted process and the alternative sequential process. Revealing in detail the influence of solvent–solute interactions remains the subject of future studies, which could also explore whether the structure of the solute before dissociation<sup>20</sup> influences the excited-state branching ratio between the different pathways.

We find that the ligation capability of  $\text{Fe}(\text{CO})_4$  is determined mostly by its  $d_{\sigma}^*$  LUMO, which receives  $\sigma$  donation from occupied CO or EtOH ligand orbitals. Population of the antibonding  $d_{\sigma}^*$  orbital in excited singlet ( $^1\text{B}_2$ ) and triplet ( $^3\text{B}_2$ )  $\text{Fe}(\text{CO})_4$  impedes  $\sigma$  donation from ligands (see the sketches in Fig. 3), explaining the inertness of these species against ligation; this problem is absent in the ligation channel that produces coordinatively saturated species. Establishment of this correlation of orbital symmetry with spin multiplicity and reactivity<sup>27</sup> is enabled by the atom specificity with which X-ray laser-based femtosecond-resolution spectroscopy can explore frontier-orbital interactions. This ability gives unique access to the reaction mechanisms of metal complexes in a way that extends and complements methods that probe structural dynamics in ultrafast chemical processes in solution<sup>28–30</sup>.

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**Supplementary Information** is available in the online version of the paper.

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