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OPEN Critical Role of Energy Transfer **Between Terbium Ions for Suppression of Back Energy** Transfer in Nonanuclear Terbium **Clusters**

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Lanthanide (Ln(III)) complexes form an important class of highly efficient luminescent materials showing characteristic line emission after efficient light absorption by the surrounding ligands. The efficiency is however lowered by back energy transfer from Ln(III) ion to the ligands, especially at higher temperatures. Here we report a new strategy to reduce back energy transfer losses. Nonanuclear lanthanide clusters containing terbium and gadolinium ions, Tb_nGd_{9-n} clusters ([Tb_nGd_{9-n}(μ-OH)₁₀(but ylsalicylate)₁₆ 1 NO₃ $^{-}$, n = 0, 1, 2, 5, 8, 9), were synthesized to investigate the effect of energy transfer between Tb(III) ions on back energy transfer. The photophysical properties of Tb_nGd_{9-n} clusters were studied by steady-state and time-resolved spectroscopic techniques and revealed a longer emission lifetime with increasing number of Tb(III) ions in Tb_nGd_{9-n} clusters. A kinetic analysis of temperature dependence of the emission lifetime show that the energy transfer between Tb(III) ions competes with back energy transfer. The experimental results are in agreement with a theoretical rate equation model that confirms the role of energy transfer between Tb(III) ions in reducing back energy transfer losses. The results provide a new strategy in molecular design for improving the luminescence efficiency in lanthanide complexes which is important for potential applications as luminescent materials.

Energy transfer (ET) plays a major role in various molecular photofunctional systems such as fluorescent biosensors^{1–3}, photon upconverters^{4–6}, and organic light-emitting diodes (OLEDs)^{7–9}. The efficiency of ET is dependent on the distance between a donor and an acceptor with maximum effective range of 1 nm (in Dexter's mechanism)10 or 10 nm (in Förster's mechanism)11. Distance sensitivity is important in fluorescent biosensors since the efficiency of ET between a target biomolecule and a sensor changes according to the distance, allowing biological activities to be monitored through spectroscopic examination¹². Closely assembled chlorophylls in a photosystem prompt fast ET and efficiently deliver the energy obtained through absorption of sunlight to a photosynthetic center¹³⁻¹⁵. In organic photovoltaics, ET between "energy relay dyes" and sensitizing dyes has been utilized to extend the absorption region for improvement of power conversion efficiency^{16,17}. Photofunctional systems have been extensively studied in many scientific fields including chemistry, physics, and biology.

Effective use of energy transfer is also seen in luminescent lanthanide (Ln(III)) complexes, and the process is known as photosensitized energy transfer (PSET). The unique photophysical properties of Ln(III) ions include spectrally fixed, sharp, and long-lived luminescence arising from their characteristic 4f-4f transition 18-21. However, efficient luminescent materials based on Ln(III) ions require photosensitization because of the small absorption coefficient of 4f-4f transitions. Photosensitization can be realized e.g. through energy transfer from organic ligands with a large ab-sorption coefficient (acting as an antenna), leading to intense lanthanide

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luminescence^{18,22,23}. Much interest has therefore been shown in these Ln(III) complexes for their potential applications to luminescent materials such as those used in bioassays, spectral converters, and electroluminescent devices^{24–31}. PSET efficiency is qualitatively known to be high when the energy level of an excited triplet state (T_1) of a ligand is slightly higher than the (emitting) excited level of an Ln(III) ion. On the other hand, when T_1 energy is too close, the reverse process of PSET, back energy transfer (BET), occurs^{32–34}. Optimal excited state energies have been focused on in Ln(III) complexes to suppress BET and achieve high luminescence efficiency.

Green-luminescent Tb(III) complexes have been one of the most frequently reported examples of Ln(III) complexes with BET $^{29,33-37}$. Experimental studies have revealed that the energy gap between the emitting level of Tb(III) ion (5D_4 state) and T_1 state of less than $1850\,\mathrm{cm}^{-1}$ promotes fast BET and low luminescence efficiency 34,37,38 . Consequently, Tb(III) complexes with high luminescence efficiency use organic ligands with high T_1 state energy and therefore even higher singlet excited (S_1) state energy, leading to absorption range shorter than $350\,\mathrm{nm}^{34,39,40}$. Expanding this range to a longer wavelength while maintaining high luminescence efficiency leads to a larger degree of freedom in molecular design, allowing incorporation of functionality, such as sensing, chirality, and self-assembly. A radically new strategy is required to suppress BET.

Herein, we propose a new method for suppressing BET in Tb(III) complexes: ET between Tb(III) ions (TbET). Photophysical processes are based on competitive kinetics 41,42 . Since both BET and TbET from a Tb(III) ion occur from the same 5D_4 excited level, these two processes are competitive. When the contribution of TbET dominates over BET, BET is suppressed. The contribution of TbET can be increased by reducing the distance between Tb(III) ions to maximize the rate of TbET. An Ln(III) cluster would therefore be the best candidate because it consists of a core of oxygen-bridged Ln(III) ions, and hence the distance between Ln(III) ions is very short, usually about 3.7 Å $^{43-49}$. We previously reported opto-magnetic properties of a nonanuclear Tb(III) cluster 50,51 . It was found that the structural features of this cluster, especially the closely assembled Tb(III) ions and the strong interaction between them, enhances the opto-magnetic property. The cluster is an ideal model to investigate the relationship between BET and TbET and to demonstrate the role of TbET in suppressing BET.

In this study, we introduced Gd(III) ions to nonanuclear Tb(III) clusters in order to investigate the relationship between TbET and luminescence efficiency (Fig. 1a). The energy level of an excited state of Gd(III) ions is high (>31500 cm⁻¹), allowing Gd(III) ions to act as spacers between Tb(III) ions to modulate the TbET processes in otherwise identical clusters. A detailed theory on how TbET can lead to suppression of BET in a nonanuclear Tb(III) cluster model is given in the "Theoretical Background" section. Novel nonanuclear Ln(III) clusters containing Tb(III) and Gd(III) [Tb_nGd_{9-n}(μ -OH)₁₀(butylsalicylate)₁₆]+NO₃⁻ (n = 0, 1, 2, 5, 8, 9) (referred to as "Tb_nGd_{9-n} clusters") were synthesized by complexation reaction of Tb(III)/Gd(III) nitrate salt and butylsalicylate. Their characteristic structures were evaluated using FAB-MS, XRD, and X-ray single crystal analysis. Emission spectroscopy, determination of quantum yield, and emission lifetime measurements were performed for the clusters in a chloroform solution. Temperature dependency of emission lifetimes was also determined for estimation of the activation energy of BET. The fundamental photophysics of Ln(III) clusters and their potential as new class of highly efficient luminescent materials are presented in this report.

Results and Discussion

Theoretical Background. The theoretical background of the concept that back energy transfer (BET) can be suppressed using energy transfer between Tb(III) ions (TbET) in a nonanuclear Tb(III) cluster is provided in this section by using rate equations. The kinetic energy transfer processes in Tb₉ clusters⁵⁰ is depicted in Fig. 1b-h. The energy transfer processes in the cluster include photosensitized energy transfer (PSET), BET, and TbET. Although the closest Tb(III)-Tb(III) pair is less than 4 Å where multipolar and exchange interactions may contribute to the energy transfer rate⁵², calculation of this contribution does not significantly affect the results and conclusion of this section (see Supplementary Information S2 and Table S1). Therefore for the sake of discussion, the energy transfer rate constant between the closest pair of Tb(III) ions was defined as k_{TbET} and the transfer rates for the other pairs were calculated relative to k_{TDET} based on the simplest R^{-6} distance dependence of the dipole-dipole Förster mechanism. Tb(III) ions in the outer unit (Tbm, m = 1, 2, ..., 8) are involved in PSET and BET with their coordinated butylsalicylate ligands. On the other hand, the center Tb(III) ion (Tb9) is coordinated only by oxygen atoms and is isolated from PSET and BET. The population density of Tb9 is solely dependent on the cycle of repetitive TbET to and from Tb(III) ions in the outer unit Tbm (energy migration). Based on these considerations, a system of differential equation in matrix form (Equation (1)) models the excited state dynamics of Tb₉ clusters:

$$\frac{\mathrm{d}X(t)}{\mathrm{d}t} = AX(t) + J(t),\tag{1}$$

where X(t) is the vector of population density of an excited singlet state (S1(t)), excited triplet state (T1(t)), excited Tb(III) ions in the outer unit (Tbm(t), m = 1, 2, ..., 8), and the center Tb(III) ion (Tb9(t)) as shown in Equation (2)

$$X(t) = \begin{pmatrix} S1(t) \\ T1(t) \\ Tb1(t) \\ \vdots \\ Tb8(t) \\ Tb9(t) \end{pmatrix}, \tag{2}$$

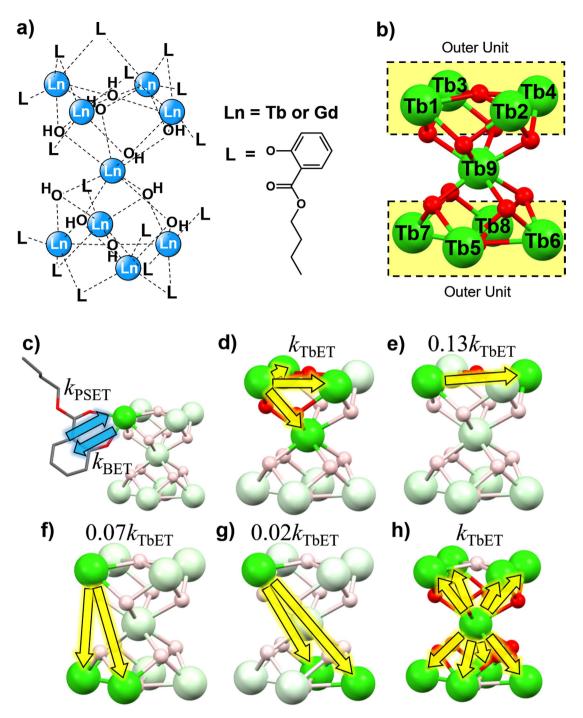


Figure 1. (a) Nonanuclear Tb(III) cluster (Tb₉ cluster). (b) Labeling of Tb₉ cluster. The center Tb(III) ion is denoted Tb9, and the Tb(III) ions in the outer unit are denoted Tbm (m = 1-8). (c) Energy transfer between Tbm and butylsalicylate ligand (only one ligand is depicted for clarity). (d) Energy transfer between Tb(III) ions from the point of view of Tb1 separated by 3.65 Å, (e) by 5.10 Å, (f) 5.64 Å, and (g) 7.10 Å. The same processes apply for all other Tbm. (h) Energy transfer between Tb(III) ions from the point of view of Tb9. The energy transfer rate constant is defined as $k_{\rm TbET}$ for two Tb(III) ions separated by 3.65 Å, and other Tb(III) ion pairs are defined relative to this rate constant according to Förster's mechanism.

A is the matrix of constants that characterizes the relationship between dynamics of each species and other species in Tb₉ cluster and is given in the following equation:

where

$$\frac{1}{\tau_{\rm S1}} = k \mathbf{r}_{\rm S1} + k \mathbf{n} \mathbf{r}_{\rm S1} + k_{\rm isc},\tag{4}$$

$$\frac{1}{\tau_{\text{T1}}} = k \mathbf{r}_{\text{T1}} + k \mathbf{n} \mathbf{r}_{\text{T1}} + 8 k_{\text{PSET}}, \tag{5}$$

$$\frac{1}{\tau_{\text{Tbm}}} = k r_{\text{Tb}} + k n r_{\text{Tb}} + 3.31 k_{\text{TbET}} + k_{\text{BET}}, \tag{6}$$

$$\frac{1}{\tau_{\text{Tb9}}} = k r_{\text{Tb}} + k n r_{\text{Tb}} + 8 k_{\text{TbET}}. \tag{7}$$

 $k_{\rm isc}$, kr, knr, $k_{\rm PSET}$, $k_{\rm BET}$, and $k_{\rm TDET}$ are defined as intersystem-crossing, radiative, nonradiative, PSET, BET, and TbET rate constants, respectively. The subscript indicates a species in a Tb₉ cluster. Finally, vector J(t) represents the input function of each species in Tb₉ cluster:

$$J(t) = \begin{pmatrix} J_{S1}(t) \\ J_{T1}(t) \\ J_{Tb1}(t) \\ \vdots \\ J_{Tb8}(t) \\ J_{Tb9}(t) \end{pmatrix}.$$
(8)

The excited state decay dynamics after short-pulse excitation were solved using Dirac's delta function, and the quantum yield (steady-state excitation) was determined by using constants (Supplementary Information S3). In this calculation, the ground state population density is approximated as always being constant, and doubly excited Tb₉ clusters are ignored because excitation with reasonably low pump intensity (under 20 W mm^{-2}) is being considered⁴¹. The population density of S1(t) (excited singlet state) after excitation is normalized to unity.

Figure 2 shows the calculated decay curve of the population density of Tb₉ cluster following a short-pulse excitation for $k_{\rm TbET}=0\,{\rm s}^{-1}$ (absence of TbET: blue line) or $k_{\rm TbET}=50000\,{\rm s}^{-1}$ (presence of TbET: red line) in the presence of BET. The other rates $k_{\rm TS1}+kn_{\rm TS1}=3.4\times10^8\,{\rm s}^{-1}$, $ki{\rm sc}_{\rm s1}=7.5\times10^8\,{\rm s}^{-1}$, $kr_{\rm T1}+knr_{\rm T1}=15000\,{\rm s}^{-1}$, and $kr_{\rm Tb}=knr_{\rm Tb}=400\,{\rm s}^{-1}$ were chosen from previous reports^{53,54}. The values of $k_{\rm PSET}$ and $k_{\rm BET}$ are arbitrary and do not affect the qualitative result. For clarity of the results, $k_{\rm PSET}=6000\,{\rm s}^{-1}$ and $k_{\rm BET}=3000\,{\rm s}^{-1}$ were chosen. The lifetime of collective Tb(III) ions (Tb1(t) + Tb2(t) + ... + Tb9(t)) is extended to $\tau_{\rm calc}=720\,{\rm \mu s}$ in the presence of TbET compared to $\tau_{\rm calc}=685\,{\rm \mu s}$ in the absence of TbET. The decay curves of the T₁ state in Fig. 2 contain short (0–0.1 ms region) and long components (0.1 ms and onward). The short component corresponds to the relaxation process of the T₁ state following intersystem crossing from the S₁ state. The time of this process matches that of the rise in population density of Tb(III) ions. The longer component corresponds to the repopulation of the T₁ state by BET from Tb(III) ions and as a result the decay rate is the same as for the Tb(III) ions. Population density of

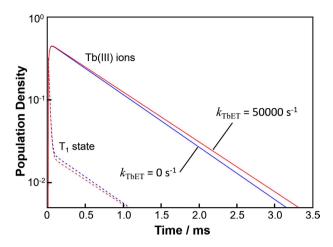


Figure 2. Calculated time evolution of population density of Tb₉ cluster after short-pulse excitation of the ligands when TbET is present ($k_{\text{TbET}} = 50000 \, \text{s}^{-1}$, red) and when TbET is absent ($k_{\text{TbET}} = 0 \, \text{s}^{-1}$, blue). The solid and dashed lines represent the collective population density of Tb(III) ions and T₁ state, respectively.

$k_{ m TbET}/{ m s}^{-1}$	$ au_{ m calc}/\mu$ s	$oldsymbol{\Phi}_{\pi\pi^*,\mathrm{calc}}$	$\eta_{ ext{BET,calc}}$
50000 (presence of TbET)	720	14.1%	44.3%
0 (absence of TbET)	685	13.4%	47.2%

Table 1. Calculated lifetimes, quantum yields and BET efficiencies of Tb₂ cluster.

the longer component of the T_1 state in the presence of TbET is lower than that in the absence of TbET, indicating suppression of BET. Table 1 summarizes the emission quantum yield $\Phi_{\pi\pi^*,\text{calc}}$ and BET efficiency $\eta_{\text{BET,calc}}$ obtained by the calculation for steady-state excitation. BET efficiency $\eta_{\text{BET,calc}}$ is defined as the difference in yield of excited Tb(III) ions with BET ($k_{\text{BET}} = 3000 \, \text{s}^{-1}$) and without BET ($k_{\text{BET}} = 0 \, \text{s}^{-1}$). An approximately 3% decrease in BET efficiency $\eta_{\text{BET,calc}}$ is observed in the Tb₉ cluster in the presence of TbET compared to that in the absence of TbET. We also calculated the population density for a hypothetical cluster where all Tb(III) ions are in an identical environment to confirm the importance of Tb9 being isolated from BET (Supplementary Information S5 and Table S3). In this case, no change was observed in the emission lifetime, quantum yield, or BET efficiency regardless of the rate of TbET. These calculations clearly show that the combination of fast TbET and Tb9 plays a role in the suppression of BET in Tb₉ cluster.

Structure and Identification. X-ray single-crystal analysis of Gd_9 cluster was performed in order to understand the molecular structure and coordination geometry of the nine Ln(III) ions in Tb_nGd_{9-n} clusters. The analysis showed that Gd(III) ions in the nonanuclear cluster take the form of an "hour-glass" structure in which the upper four and lower four Gd(III) ions are connected to the center Gd(III) ion via oxygen-bridging as shown in Fig. 3a. All Gd(III) ions take the form of an 8-coordination structure. The center Gd(III) ion is coordinated only by eight oxygen atoms. Each of the Gd(III) ions in the upper four and lower four sites (outer unit) are coordinated by both butylsalicylate ligands and oxygen atoms. The distances between two Gd(III) ions for all combinations are summarized in Table S4.

Evaluation of coordination geometry of Gd(III) ions was done through continuous shape measures (CShM) calculation using SHAPE^{55–57}. The CShM criterion *S*, summarized in Table S5, represents the degree of deviation from ideal coordination geometry. We have chosen CShM over another method called "shape measure" (ShM)⁵⁸ because it takes account of the distortion of center metal ion from the center of mass. The principle and comparison of the calculation method in CShM and ShM is described in Supplementary Information S7. The *S* value for the center Gd(III) ion (Fig. 3b) was 0.082 when calculated for 8-coordinated square antiprism (8-SAP) geometry and it was 2.481 for 8-coordinated trigonal dodecahedron (8-TDH) geometry. Therefore, the coordination geometry of the center Gd(III) ion is attributed to the 8-SAP structure. The upper four and lower four Gd(III) ionns (Fig. 3c) take an 8-TDH structure since the *S* values calculated for 8-TDH are smaller than those for 8-SAP. This structural result is similar to that of our previously reported Tb₉ cluster⁵⁰.

The structures of the other Tb_nGd_{9-n} clusters (n=1,2,5,8) were determined by the combination of powder XRD (Figure S2a) and FAB-MS (Figure S2b). In the XRD results, the diffraction angle is corrected on the Si peak at $2\theta=28.43$ degrees for accurate comparison of the peaks derived from the Tb_nGd_{9-n} clusters. Three distinguished peaks ($2\theta_1=11.3$ degrees, $2\theta_2=15.6$ degrees, and $2\theta_3=19.8$ degrees) were observed for all of the clusters indicating that the mixed Tb/Gd clusters have the same structure as those of Tb_9 and Tb_9 and Tb_9 and Tb_9 clusters. This is further supported by FAB-MS spectroscopy which showed Tb_9 all of the Tb_9 clusters corresponding to their calculated molecular weight without the Tb_9 counter-anion. The results presented above indicate that all Tb_9 Tb_9 clusters were successfully synthesized with identical structures.

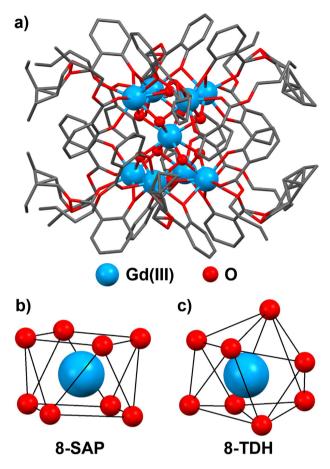


Figure 3. Molecular structure of Gd₉ cluster. (a) Structure of Gd₉ cluster (hydrogen atoms and NO₃⁻ counter anion omitted for clarity), (b) Gd(III) ion with 8-SAP geometry, and (c) Gd(III) ion with 8-TDH geometry.

Optical Properties. Emission spectra of $\mathrm{Tb}_n\mathrm{Gd}_{9-n}$ clusters (n=1,2,5,8,9) in $1.0\times10^{-4}\mathrm{M}$ chloroform solution were measured and Fig. 4a shows the spectra, normalized at the peak intensity. The Tb_nGd_{9-n} clusters exhibit characteristic emission of the 4f-4f transition of the Tb(III) ion with each of the peaks corresponding to its ${}^5D_4 \rightarrow {}^7F_I(J=6-1)$ transitions 59 . The spectral shape of the emission was identical in all Tb_nGd_{9-n} clusters. Additionally, the emission spectral shape of Tb₉ cluster in powder form (Figure S4) was the same as that in the solution form, indicating that the coordination structure is maintained in solution. The emission spectrum of Gd_9 cluster in $1.0 \times 10^{-4} M$ chloroform solution at 210 K (Figure S5) was measured to estimate the excited triplet state (T_1) energy level of the ligand ^{38,60,61}. The emission peak was observed at 461 nm (21690 cm⁻¹). Since the 5D_4 level of a Tb(III) ion is $20620 \,\mathrm{cm}^{-1}$ as observed in the emission spectra (485 nm), the energy gap ΔE between T₁ and ⁵D₄ was estimated to be 1070 cm⁻¹. This is well within the range for BET in Tb(III) complexes $(\Delta E < 1850 \, \text{cm}^{-1})^{33,34,38}$, and thus BET is expected to take place in Tb_nGd_{9-n} clusters. The absolute emission quantum yield was determined by using an integration sphere ($\lambda_{\rm EX} = 380$ nm, ligand excitation) and emission lifetimes were determined by using a nanosecond pulse laser ($\lambda_{EX} = 355 \text{ nm}$, $\lambda_{EM} = 550 \text{ nm}$, ligand excitation) for clusters in a chloroform solution. The emission decay profiles were single-exponential (Fig. 4b) in all Tb_nGd_{9-n} clusters. As summarized in Table 2, the emission quantum yield $\Phi_{\pi\pi^*}$ and lifetime $au_{
m obs}$ increased with increasing numbers of Tb(III) ions in the Tb_nGd_{9-n} clusters. The emission quantum yield reflects the effect of all photophysical processes in Tb_nGd_{0-n} clusters. Of these processes, TbET, PSET, and BET are all likely to vary with n in the Tb_nGd_{0-n} clusters, and the effect purely based on TbET is indistinguishable.

The effect of TbET with increasing number of Tb(III) ions in the clusters was investigated by measuring the temperature dependence of emission lifetimes (Fig. 5). A clear temperature dependence was observed above $270 \, \text{K}$ for $\text{Tb}_1 \text{Gd}_8$ and $\text{Tb}_2 \text{Gd}_7$ clusters, while lifetimes remained constant for $\text{Tb}_5 \text{Gd}_4$, $\text{Tb}_8 \text{Gd}_1$, and Tb_9 clusters. The rate constant of BET k_{BET} can be estimated by the following equation 37,62 .

$$\ln k_{\rm BET} = \ln \left(\frac{1}{\tau} - \frac{1}{\tau_{210\rm K}} \right) = \ln A - \frac{E a_{\rm BET}}{RT},\tag{8}$$

where τ is measured emission lifetime at a given temperature, $\tau_{210\text{K}}$ is the emission lifetime at 210 K, A is frequency factor, Ea_{BET} is activation energy for BET, R is gas constant, and T is temperature. Ea_{BET} and k_{BET} can be calculated for Tb₁Gd₈ cluster by an Arrhenius plot within the range in which temperature dependency was

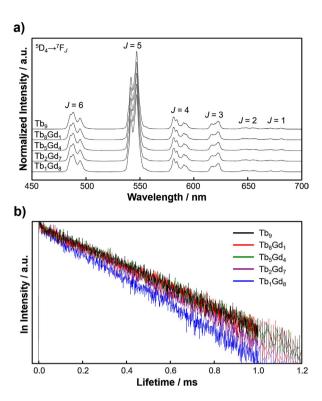


Figure 4. (a) Photoluminescence spectra of Tb_nGd_{9-n} clusters in $1.0\times 10^{-4}\,\mathrm{M}$ chloroform solution normalized at the peak top. Excitation wavelength λ_{EX} was 380 nm. (b) Normalized emission lifetimes of Tb_nGd_{9-n} clusters in $1.0\times 10^{-4}\,\mathrm{M}$ chloroform solution at 300 K. All spectra were normalized at emission intensity at 0 seconds. All samples were excited by nanosecond pulse laser at $\lambda_{EX} = 355\,\mathrm{nm}$ and monitored at $\lambda_{EM} = 550\,\mathrm{nm}$ ($^5D_4 \rightarrow ^7F_6$ transition).

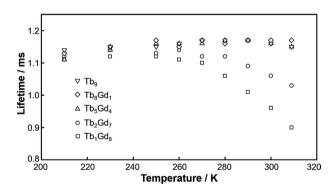


Figure 5. Emission lifetime temperature dependency of Tb_nGd_{9-n} clusters in $1.0 \times 10^{-4} M$ chloroform solution. The experimental uncertainty in the reported lifetimes is up to $\pm 1.5\%$.

Clusters	$oldsymbol{\Phi}_{\pi\pi^*,}{}^a$	$ au_{ m obs}/ m ms^{ m b}$	$ au_{210 ext{K}}/ ext{ms}^ ext{b}$	$A/10^8 \mathrm{s}^{-1c}$	$k_{\rm BET}/{\rm s}^{-1{\rm c}}$	Ea _{BET} /kj mol ^{-1d}
Tb ₁ Gd ₈	14%	0.96	1.11	9.10	167	38.7
Tb ₂ Gd ₇	23%	1.06	1.12	*	*	*
Tb ₅ Gd ₄	33%	1.16	1.11	*	*	*
Tb ₈ Gd ₁	40%	1.17	1.13	*	*	*
Tb ₉	39%	1.16	1.14	*	*	*

Table 2. Photophysical properties of Tb_nGd_{9-n} clusters. ^aMeasured in 1.0×10^{-4} M chloroform solution ($\lambda_{\rm EX} = 380$ nm). ^bMeasured in 1.0×10^{-4} M chloroform solution ($\lambda_{\rm EX} = 355$ nm). ^cCalculated from Equation (9). A is frequency factor. ^dAnalyzed from an Arrhenius plot of Equation (9) using lifetime temperature dependency results. $k_{\rm BET}$ and $E_{a\rm BET}$ could only be calculated for Tb₁Gd₈ cluster since other Tb_nGd_{9-n} clusters (n = 2, 5, 8, 9) involve TbET, which contributes to the temperature dependency of lifetimes.

observed (Figure S6). The BET rate constant $k_{\rm BET}$ and activation energy $Ea_{\rm BET}$ for Tb₁Gd₈ cluster were calculated to be $167~{\rm s}^{-1}$ and $38.7~{\rm kJ~mol}^{-1}$, respectively (Table 2).

Emission lifetime and its temperature dependence are in qualitative agreement with theoretical results presented in the Theoretical Background section. The theoretical lifetime of Tb $_9$ cluster in the absence of BET is constant regardless of TbET rate $k_{\rm TbET}$ (Supplementary Information S4 and Table S2). Experimentally, BET is absent in Tb $_{\rm n}$ Gd $_{\rm 9-n}$ clusters at 210 K, and the number of Tb(III) ions in the cluster is directly related to $k_{\rm TbET}$. Thus the emission lifetime at this temperature should be the same for all Tb $_{\rm n}$ Gd $_{\rm 9-n}$ clusters. The experimental emission lifetimes of Tb $_{\rm n}$ Gd $_{\rm 9-n}$ clusters at 210 K $\tau_{\rm 210K}$ are approximately 1.12 ms regardless of the number of Tb(III) ions in the clusters (Table 2 and Fig. 5). Meanwhile, the theoretical lifetime of Tb $_{\rm 9}$ clusters in the presence of BET shows that the lifetime is longer in the presence of TbET than in the absence of TbET. This comparison is analogous to experimental lifetimes of Tb $_{\rm 9}$ cluster (presence of TbET) and Tb $_{\rm 1}$ Gd $_{\rm 8}$ cluster (absence of TbET) at a temperature above 270 K, at which BET occurs. At a temperature above 270 K, the experimental lifetime is clearly longer for clusters with larger numbers of Tb(III) ions than for clusters with smaller numbers of Tb(III) ions. Such agreement between the theoretical and experimental results implies that the contribution of BET is indeed suppressed in Tb $_{\rm 9}$ cluster because of TbET and the existence of a center Tb(III) ion not in direct contact with ligands for which no BET occurs.

Summary and Conclusion. The role of energy transfer between Tb(III) ions in luminescent Tb(III) complexes has been demonstrated to reduce BET losses both theoretically and experimentally using Tb_nGd_{9-n} clusters ($[Tb_nGd_{9-n}(\mu\text{-OH})_{10}(\text{butylsalicylate})_{16}]^+NO_3^-$) as model system. In the Theoretical Background section, we demonstrated that two features of Tb₉ cluster were important for suppressing BET: 1) TbET by closely assembled Tb(III) ions and 2) existence of Tb9 that is isolated from BET. Experimentally, we synthesized and investigated the photophysical properties of nonanuclear Tb_nGd_{9-n} clusters where n=0,1,2,5,8, and 9. The combination of X-ray single-crystal analysis, powder XRD, and FAB-MS revealed that these clusters have nearly identical structures. Temperature dependency measurements of emission lifetime revealed that effect of BET becomes prominent in Tb_nGd_{9-n} clusters at temperatures above 270 K. Below this temperature, the emission lifetime is constant for all clusters. Above 270 K, the decrease in emission lifetime is mitigated for clusters with over five Tb(III) ions. By comparing the trends observed in experimental results to those of theoretical results, it was found that the contribution of BET is indeed suppressed in clusters with a large number of Tb(III) ions. These findings provide a new insight into the fundamental photophysics of Ln(III) complexes as well as indication of a novel strategy to achieve higher luminescence efficiency in Ln(III) complexes.

Methods

Materials. Terbium(III) nitrate hexahydrate (Tb(NO₃)₃·6H₂O) and triethylamine ($C_6H_{15}N$) were purchased from Kanto Chemical Co. Inc. Gadolinium(III) nitrate hexahydrate (Gd(NO₃)₃·6H₂O) was purchased from Sigma-Aldrich Co. Butylsalicylate ($C_{11}H_{14}O_3$) was purchased from Tokyo Chemical Industry Co., LTD. All other chemicals and solvents were of reagent grade and were used without further purification.

Apparatus. FAB-MS spectra were measured on a JEOL JMS-700TZ. Elemental analyses were performed by Exter Analytical CE440. Infrared spectra were recorded on a JASCO FT/IR-4600 spectrometer. XRD spectra were characterized by a RIGAKU X-ray diffractometer RINT 2200. Single crystal X-ray diffractions were made on a RIGAKU RAXIS RAPID imaging plate area detector.

Optical Measurements. Absorption spectra of Tb_nGd_{9-n} clusters were obtained by using a JASCO V-670 spectrometer. Emission spectra were measured using a Horiba/Jobin-Yvon FluoroLog-3 spectrofluorometer and a JASCO FP-6600 spectrometer. The combination of an integration sphere and a JASCO FP-6600 spectrometer was used to measure emission quantum yields. Emission lifetimes were measured using the third harmonic (355 nm) of a Q-switched Nd:YAG laser. The temperature was controlled using an Oxford Instruments OptistatDN2 cryostat.

Crystallography. Colorless single crystals of Gd_9 cluster obtained from solutions in methanol were mounted on a glass fiber by using epoxy resin glue. All measurements were made using a Rigaku RAXIS RAPID imaging plate area detector with graphite-monochromated $MoK\alpha$ radiation. Corrections for decay and Lorentz-polarization effects were made using a spherical absorption correction, solved by direct methods, and expanded using Fourier techniques. Non-hydrogen atoms were refined anisotropically except for disordered atoms. Hydrogen atoms were refined using the riding model. The final cycle of full-matrix least-squares refinement was based on observed reflections and variable parameters. All calculations were performed using a CrystalStructure crystallographic software package. We confirmed the CIF data by using the checkCIF/PLATON service. CCDC-1479981 (Gd_9 cluster) contains the crystallographic data for this paper. These data can be obtained free of charge from The Cambridge Crystallographic Data Center via www.ccdc.cam.ac.uk/data_request/cif.

Synthesis of Gd₉ cluster ([Gd₉(\mu-OH)₁₀(butylsalicylate)₁₆]⁺NO₃⁻). Butylsalicylate (1.05 g, 5.40 mmol) was dissolved in methanol, and triethylamine (1.22 ml, 8.80 mmol) was added to the solution and the solution was stirred at 40 °C. Then Gd(NO₃)₃·6H₂O (1.372 g, 3.04 mmol) in methanol was added dropwise to the solution and the solution stirred for 30 minutes. The solution was cooled to room temperature, and white powder Gd₉ cluster ([Gd₉(μ -OH)₁₀(butylsalicylate)₁₆]⁺NO₃⁻) was obtained 15 minutes later. Yield: 59%. Selected IR (ATR, cm⁻¹): 3573 (w, O-H), 3234 (w, O-H), 2956 (m, C-H), 1318 (s, C-O). FAB-MS: m/z = 4676.8 [Gd₉(μ -OH)₁₀(butylsalicylate)₁₆]⁺. Elemental analysis calculated for C₁₇₆H₂₁₈NO₆₁Gd₉: C, 44.61%, H, 4.64%, N, 0.30%. Found: C, 44.30%, H, 4.53%, N, 0.25%.

Synthesis of Tb₁Gd₈ cluster ([Tb₁Gd₈(\mu-OH)₁₀(butylsalicylate)₁₆]⁺NO₃⁻). Tb₁Gd₈ cluster was synthesized by the same procedure as that used for Gd₉ cluster except for the use of Tb(NO₃)₃·6H₂O (0.172 g, 0.380 mmol) and Gd(NO₃)₃·6H₂O (1.24 g, 2.73 mmol). Yield: 82%. Selected IR (ATR, cm⁻¹): 3573 (w, O-H), 3234 (w, O-H), 2956 (m, C-H), 1318 (s, C-O). FAB-MS: m/z = 4677.8 [Tb₁Gd₈(μ -OH)₁₀(butylsalicylate)₁₆]⁺. Elemental analysis calculated for C₁₇₆H₂₁₈NO₆₁Tb₁Gd₈: C, 44.59%, H, 4.64%, N, 0.30%. Found: C, 43.95%, H, 4.53%, N, 0.29%.

Synthesis of Tb₂Gd₇ cluster ([Tb₂Gd₇(μ-OH)₁₀(butylsalicylate)₁₆]⁺NO₃⁻). Tb₂Gd₇ cluster was synthesized by the same procedure as that used for Gd₉ cluster except for the use of Tb(NO₃)₃·6H₂O (0.306 g, 0.675 mmol) and Gd(NO₃)₃·6H₂O (1.07 g, 2.36 mmol). Yield: 51%. Selected IR (ATR, cm⁻¹): 3573 (w, O-H), 3234 (w, O-H), 2956 (m, C-H), 1318 (s, C-O). FAB-MS: m/z = 4679.8 [Tb₂Gd₇(μ-OH)₁₀(butylsalicylate)₁₆]⁺. Elemental analysis calculated for C₁₇₆H₂₁₈NO₆₁Tb₂Gd₇: C, 44.58%, H, 4.63%, N, 0.30%. Found: C, 44.32%, H, 4.54%, N, 0.26%.

Synthesis of Tb₅Gd₄ cluster ([Tb₅Gd₄(\mu-OH)₁₀(butylsalicylate)₁₆]⁺NO₃⁻). Tb₅Gd₄ cluster was synthesized by the same procedure as that used for Gd₉ cluster except for the use of Tb(NO₃)₃·6H₂O (0.688 g, 1.52 mmol) and Gd(NO₃)₃·6H₂O (0.684 g, 1.52 mmol). Yield: 41%. Selected IR (ATR, cm⁻¹): 3573 (w, O-H), 3234 (w, O-H), 2956 (m, C-H), 1318 (s, C-O). FAB-MS: m/z = 4684.9 [Tb₅Gd₄(μ -OH)₁₀(butylsalicylate)₁₆]⁺. Elemental analysis calculated for C₁₇₆H₂₁₈NO₆₁Tb₅Gd₄: C, 44.53%, H, 4.63%, N, 0.30%. Found: C, 44.23%, H, 4.52%, N, 0.24%.

Synthesis of Tb₈**Gd**₁ **cluster ([Tb**₈**Gd**₁(μ -**OH)**₁₀(**butylsalicylate)**₁₆]⁺**NO**₃⁻). Tb₈**Gd**₁ cluster was synthesized by the same procedure as that used for **Gd**₉ cluster except for the use of Tb(NO₃)₃·6H₂O (1.026 g, 2.26 mmol) and **Gd**(NO₃)₃·6H₂O (0.344 g, 0.762 mmol). Yield: 49%. Selected IR (ATR, cm⁻¹): 3573 (w, O-H), 3234 (w, O-H), 2956 (m, C-H), 1318 (s, C-O). FAB-MS: m/z = 4689.8 [Tb₈**Gd**₁(μ -**OH**)₁₀(butylsalicylate)₁₆]⁺. Elemental analysis calculated for C₁₇₆H₂₁₈NO₆₁Tb₈**Gd**₁: C, 44.48%, H, 4.62%, N, 0.29%. Found: C, 44.24%, H, 4.53%, N, 0.24%.

Synthesis of Tb₉ **cluster ([Tb**₉(μ -OH)₁₀(**butylsalicylate**)₁₆]⁺NO₃⁻). Tb₉ cluster was synthesized by the same procedure as that used for Gd₉ cluster except for the use of Tb(NO₃)₃·6H₂O (1.372 g, 3.04 mmol). Yield: 52%. Selected IR (ATR, cm⁻¹): 3573 (w, O-H), 3234 (w, O-H), 2956 (m, C-H), 1318 (s, C-O). FAB-MS: m/z = 4691.3 [Tb₉(μ -OH)₁₀(butylsalicylate)₁₆]⁺. Elemental analysis calculated for C₁₇₆H₂₁₈NO₆₁Tb₉: C, 44.47%, H, 4.62%, N, 0.29%. Found: C, 44.20%, H, 4.57%, N, 0.27%.

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Author Contributions

S.O. performed all the syntheses and measurements, and wrote the paper under the supervision of T.N. and Y.H. S.O., T.N., Y.K., K.F., A.M., and Y.H. discussed and designed the research. T.S. and H.I. supported the single crystal and power XRD measurements. All authors reviewed the manuscript.

Additional Information

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